

# Time Transient Analysis and Non-Linear Rotordynamics

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Abstract: Rotordynamics analysis is used to design and understand the vibration of rotating equipment. Fortunately, the vast majority of rotordynamics analyses are done in the frequency domain and are linear in nature. In reality, almost nothing is purely linear, but the non-linear aspects are so small that they can usually be safely ignored. The normal analysis techniques are useful and accurate enough to predict critical speeds, vibration amplitudes in response to imbalances, misalignments, bowed rotors, vane-pass excitations and other common harmonic forcing functions. When calculating undamped lateral or torsional critical speeds and mode shapes only linear springs and masses are used. When we conduct a damped analysis like forced synchronous response or stability, factors that are somewhat nonlinear like the stiffness and damping of journal bearings are essentially linearized by calculating discrete values at each speed considered. If we put in one unit of imbalance, we get out a certain amount of predicted vibration amplitude at each rotor station and at each speed. If we double the amount of imbalance, the output is exactly doubled.

There are, however, many cases where the frequency domain does not provide adequate analysis capability. These are the times when time transient analysis becomes necessary. We need to do this because the forcing functions vary with time as with the pulsations from a synchronous motor during a startup. The system may be subjected to a complex forcing function with many harmonics like a reciprocating compressor or the system may be subjected to an sudden change like a complete generator load shed or the loss of a turbine blade. In all these cases the steady state response is altered by the non-linear input and the system response must be calculated as a function of time. Many times we are concerned with component stresses and fatigue life. One example in this paper demonstrates how the fatigue life can be calculated from a time based output.

The other primary non-linear phenomena that occurs is when a component in the system does not act linearly as a function of amplitude. Some examples are hydrodynamic bearings, squeeze-film damper bearings, elastomeric couplings, magnetic bearings, electrical fields, and support structures. This paper will consider high amplitudes in plain journal bearings and squeeze-film damper bearings. Anyone who has performed a field balance on a fan with large vibration will have experienced the non-linear effect. The initial trial weight prediction will usually be incorrect until the vibration is reduced into a more linear response range.

**Key Words:** Amplitude, Bearings, Blade Loss, Compressor, Dampers, Damping, Fatigue, Gears, Lateral, Linear, Load Shed, Motor, Orbits, Non-linear, Reciprocating, Squeeze-Film, Stress, Synchronous Motor, Time, Torsional, Turbine, Vibration

**Introduction:** Traditional lateral and torsional rotordynamics analyses involve taking the physical dimensions of a rotor system and its associated bearings and supports and building a mathematical analog. Many papers by this author and books and papers by others listed in the references describe in detail these methods, capabilities and limitations. The traditional analysis approach views the dynamics of a rotor-bearing system as a function of frequency. In this paper we will explore rotordynamics in the time domain and, by extension into the non-linear regime.

For train torsional analysis we are primarily interested in resonant frequencies, mode shapes and forced response due to torque pulsations. While standard undamped torsional resonant frequencies and mode shapes are very easy to calculate, one must be very careful to define the alternating applied torques when attempting to calculate torsional twist amplitudes with a forced response calculation. Alternating torques may be created by mechanical means from gear sets or couplings, may be a consequence of the driving torques developed by the prime mover as in a synchronous motor or may be generated by the process side of the machinery like a reciprocating compressor. All of these factors become extremely important when calculating fatigue life of components that are subjected to periodic torsional shear stresses. It will be shown that one often needs to enter the time domain to adequately describe the torsional forcing functions and the resultant stresses if the goal of realistic fatigue life is to be attained.

For lateral rotordynamics we are very interested in critical speeds, mode shapes, forced response and stability. An undamped lateral critical speed analysis is very useful to give the user a feeling for the critical speeds and mode shapes. A critical speed map, especially with the predicted bearing coefficients cross-plotted, is a very useful design tool. However, for predicting actual frequencies of maximum rotor amplitudes, the vibration magnitudes that can be expected and the rotor stability it is necessary to include damping. Some of the elements in the forced response and stability analyses such as seal and bearing stiffness and damping coefficients are not linear (we calculate them as a function of rotor speed). However, in the standard linear synchronous unbalance response analysis the amplitude, phase and frequency calculations are linear at each speed step. To get the response predictions we simply connect the "dots" to produce the normal Bode type plots.

**Time Based Analysis:** There are two major elements of rotordynamics analysis that can fall into the realm of time based analysis. The first is those physical elements that cannot be defined properly as a function of frequency. Examples of this are seals, bearings, and, in particular, squeeze film dampers. The characteristics of these elements vary with amplitude or eccentricity. Sometimes the characteristics of structural elements vary with amplitude as anyone with field balancing experience can testify. Secondly, time based forcing functions must be considered with some systems. Examples include synchronous motors and reciprocating engines and compressors.

**Linear and Non-Linear:** Linear analysis can be though of as a constant multiplier. "X units of imbalance will cause "CX units of vibration amplitude as illustrated in figure 1 where "C is the constant. If you double the input magnitude, the output doubles as well. This is based on the assumption that the system elements are all linear like a simple steel spring or that the non-linearities are small (like the hysteresis in the steel) enough that they do not significantly affect the calculated results. For many rotordynamics analyses these assumptions are completely adequate and the vast majority of such work relies on this fact to produce meaningful and useful results. In the real world it is never this simple, but we're always working to close the gap between analysis and reality.

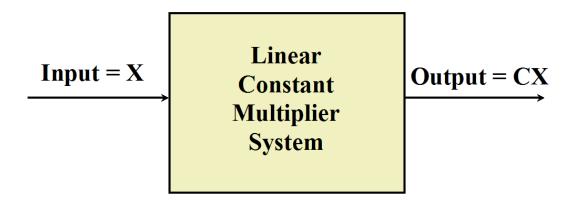


Figure 1 - Typical Linear System

Non-linear analysis is difficult because most non-linearities cannot be easily defined. Figure 2 illustrates a non-linear system. Here the input is modified as a function of time, F(t) or speed F(N), or other variable like amplitude. An offset value "Y may also be added by the system. Squeeze film dampers are a good example. The stiffness and damping characteristics of a squeeze film damper are highly dependent on just how centered the damper is in its clearance annulus and how large the vibration amplitudes are. Up to about 40 percent of the damper eccentricity, linear assumptions work fairly well but at higher eccentricities a technique must be used where the current calculated position must be used, along with the rest of the dynamic model, to calculate the next position. For a given time step, new forces and dynamic stiffness and damping coefficients are then calculated for the next position and the process repeats for each tiny time step. This way actual amplitudes and stability characteristics can be calculated for non-linear elements.

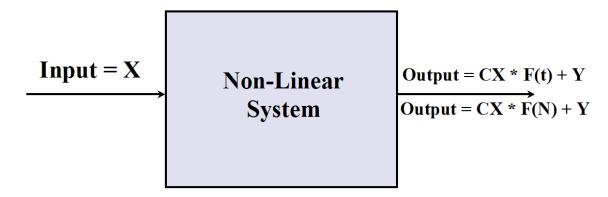


Figure 2 - Typical Non-Linear System

Non-linear input functions are also possible. The case of the driving and alternating torques from a synchronous motor during startup is a good example. These torques cannot be avoided and must be analyzed with a time transient program in order to avoid catastrophic failures. Reciprocating compressors and engines develop even and half-order harmonics that can vary with speed and various load conditions. The non-linear torques are definable as function of time. Accurate shaft fatigue life calculations depend on accurate definition of the forcing functions. Sudden step changes in torque or imbalance must also be analyzed in the time domain.

Torsional Time Transient Analysis: Two primary examples of the need for time based torsional analysis are presented here. The first is the case of synchronous motor startups. Generally this analysis is linear with respect to the assumption that all of the elements are linear springs. The nonlinear part comes in because the driving and alternating torque inputs are a function of speed. In addition, at large amplitudes there may be significant torsional-lateral interaction that can alter the amplitudes and stresses of the system. If the stresses approach yield there may be minute local deformations that relieve the highest stresses. Unfortunately, such amplitude response nonlinearities are extremely difficult to quantify and are not included in this example. Transient torsional excitation during the start of a synchronous motor can be one of the most severe cases of transient shaft shear stress and fatigue. There are cases where these systems failed catastrophically on the first startup and many more where failure occurred after multiple startups as fatigue cycles accumulated. The steel components in any system "remember every stress cycle."

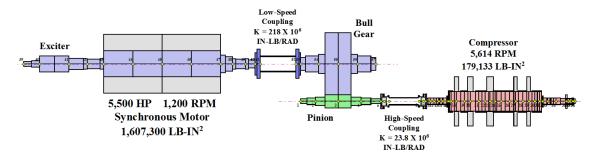


Figure 3 - Synchronous Motor-Gear-Compressor Train

Figure 3 shows a 6-pole 5,500 HP synchronous motor driving an air compressor through a gear speed increaser. Every synchronous motor generates alternating torque impulses during a startup. These pulses begin at twice line frequency (7,200 CPM in the US) and decline inversely with speed to zero when synchronous speed is reached. This means that every torsional resonance between 0 and 7,200 CPM will be excited during each startup. Thus, any machine train with a synchronous motor must be evaluated with a transient torsional analysis. A transient analysis consists of creating a springmass model and defining all the torque producers (the motor and possibly the gear) and torque absorbers (the compressor). The torque information is generally available only from the motor vendor who must communicate with the manufacturer of the other machinery and produce a set of torque curves similar to figure 4.

A forced response procedure is applied to the model simulating the startup. At time zero, the driving, oscillatory, and load torques are applied simultaneously. This causes the train to accelerate in speed. The input torques are resisted by the polar inertia of the machines and the gas load from the compressor. Resistance from the bearings and gears is neglected. Then, at very small time steps, a new speed is calculated along with the torsional oscillation amplitude at every point in the train. Ultimately, the alternating torque amplitudes can be converted to the shear stresses in the shafts. One important calculation coming from this analysis is the time to reach full speed as seen in figure 5. Typically, anything over 20 seconds could be damaging to the motor. After synchronization, the delivered motor torque is balanced by the compressor load torque and there are no more alternating torque pulsations from the motor. Very low frequency oscillations often seen immediately after synchronization are due to interaction of the mechanical system with the electrical grid.

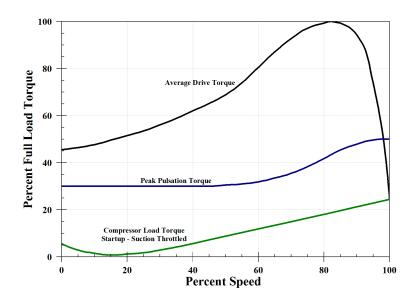


Figure 4 - Speed vs. Torque Curves for 5,500 HP Synchronous Motor

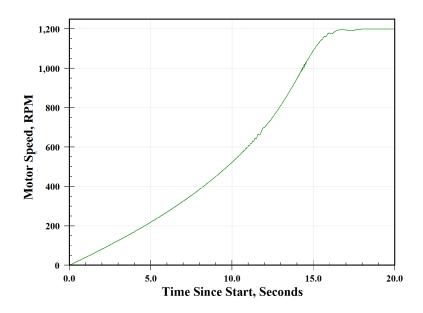


Figure 5 - Startup Time Calculation for Synchronous Motor-Gear-Compressor Train

Ripples visible in the time-speed curve of figure 5 are actually speed fluctuations due to the torsional resonances. The first disturbance, encountered 11.8 seconds into the startup, is the second torsional resonance at 3,260 CPM. The first torsional resonance, at 1,275 CPM, is encountered at 14.5 seconds into the startup. Figure 6 is the animated mode shape of the first torsional resonance. Most of the twist amplitude is between the motor and the bull gear. This is where the maximum shear stresses will occur. For this reason the coupling flanges on both the motor and the bull gear were made integral with the shafts. This avoids any stresses associated with shrunk-on hubs or keyways. It should be noted that use of non-linear elastomer (so-called "damper") couplings can alter

these startup and response curves significantly. Elastomer couplings are not covered here.

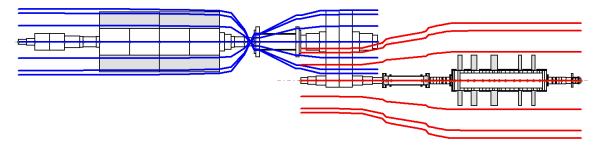


Figure 6 - First Torsional Resonance Mode Shape for Example 3 - 1,275 CPM

Figure 7 is a plot of the alternating shaft shear stresses in the motor shaft section near the coupling during a portion of the startup. The low-speed shaft always has higher torques and usually has the highest steady and alternating torques from the synchronous motor startup.

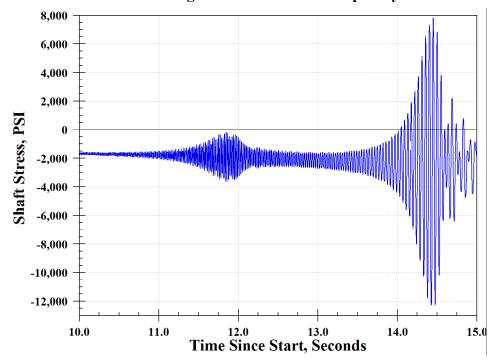


Figure 7 - Shaft Stresses During Time Transient Startup of Synchronous Motor

The results of the time transient analysis produce a history of the number and magnitude of the alternating shear stresses in the shafting. In order to determine if these fatigue cycles will ultimately be a cause of failure, a fatigue analysis is performed. It is not possible for me to do this accurately without a time transient analysis. In addition, derating factors are required such as stress concentration factors, material composition and defect factors, surface finish factors, and other things that may or may not apply to a given situation like a corrosion derating factor. Finally the time history is examined and the minimum and maximum shear stress values for the time period of high alternating stresses is examined and averaged to get a list of the maximum peak stresses.

Assuming the material specification is AISI 4140/4145 with an ultimate tensile strength (UTS) of 110,000 PSI and a yield strength of 85,000 PSI with a 40% reduction in area. A shear fatigue life plot for this material is shown in figure 8.

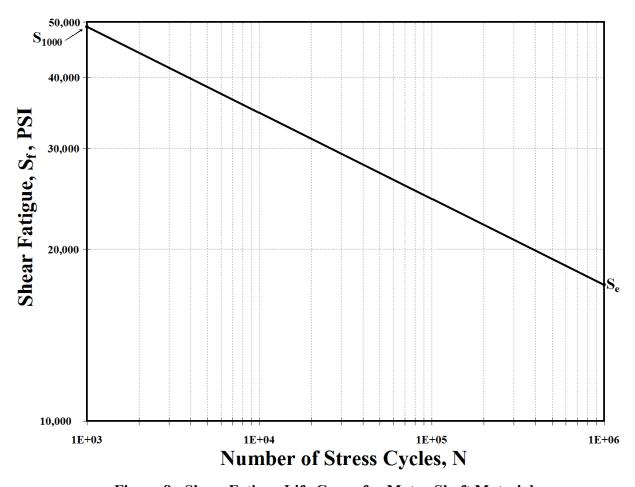


Figure 8 - Shear Fatigue Life Curve for Motor Shaft Material

In this particular analysis, failures that occur at less than 1000 cycles are called low cycle fatigue and high cycle fatigue after 1000 cycles to failure. Above 1 million cycles, the component is assumed to have infinite life. Figure 8, for torsional shear stress, was derived from a plot of tensile failure data by multiplying by 0.577 (von Mises theory) meaning that the shear endurance limit for infinite life is reduced from a tensile value of 30,000 PSI to  $S_e = 17,310$  PSI. At the other end, failure at 1000 cycles ( $S_{1000}$ ) will occur at a shear stress level of 49,000 PSI. Table 1 is a listing of the calculated shaft shear stresses found in a plot of similar to figure 7. The stress concentration factor applied is 2.5 for the keyway and a factor of 1.35 is applied as a factor of safety. Also, a surface finish derating factor of 0.9 was used. A system modal damping factor was also applied of two percent of critical damping. The choice of this factor is critical since if it is too low, every system will calculate as a failure and if too high, every system will be indicated to be acceptable. Fatigue analysis is somewhat an art and must have a basis in experience.

Cycle	Maximum Shear Stress, PSI	Minimum Shear Stress, PSI	Cyclic Stress (Max-Min)/2	With SCF=2.5, SF=.9 and FOS = 1.35
-3	7,098.3	-255.8	3,677	13,789
-2	8,642.2	-1,791.4	5,217	19,563
-1	10,725.8	-3,628.2	7,177	26,914
Peak	12,719.6	-4,457.0	8,588	32,206
1	12,490.2	-1,741.3	7,116	26,684
2	8,988.0	751.5	4,118	15,443
3	10,246.7	-467.1	5,357	20,088
4	7,758.4	1,741.6	3,008	11,281

Table 1 - Shear Stress Cycles Summary from Figure 14

Since the shear endurance limit is 17,310 PSI then only five cycles per startup will actually contribute to the fatigue of the motor shaft (these are italicized in the table). For each of these high-stress cycles we must determine what portion of the shaft life is used up. Then using Miner's rule, the cycles to failure can be calculated. Referring to figure 8, the cycles to failure for the five high stress cycles are tabulated in table 2.

Cycle	<b>Combined Shear Stress, PSI</b>	Cycles to Failure
1	19,563	3.9 X 10 <sup>5</sup>
2	26,914	8.8 X 10 <sup>4</sup>
3	32,206	1.2 X 10 <sup>4</sup>
4	26,684	8.8 X 10 <sup>4</sup>
5	20,088	3.7 X 10 <sup>5</sup>

Table 2 - Calculated Life for each Stress Cycle in each Start

Based on actual experience, most references suggest using a modification to Miner's rule of 0.5 which is, in essence, another safety factor. Thus the calculated life of the motor shaft will be:

$$N_{Starts} = \frac{0.5}{\sum \frac{1}{N_{Cycle}}}$$

The inverse sum of the inverse values in table 4 is 8,983 so the predicted number of safe startups as applied to the motor shaft is 4,491 starts. Since this is a reasonable number of safe starts, the motor shaft is acceptable.

## **Reciprocating Compressor Train Analysis**

The second case of using time transient analysis involves steady state excitation from a complicated alternating torque source. The compressor train illustrated here consists of a 600 HP induction motor driving a 4-cylinder, double acting vertical piston reciprocating compressor with two 480mm and two 330mm cylinders. The motor is directly coupled to the gear by a disk type coupling. The design full load motor speed is 593 RPM.

Figure 9 is a cross-section of the entire compressor train computer model. The provided physical dimensions and other torsional information were translated into the finite element analysis model used here. The motor rotor is shown in light red, the coupling in light blue, the flywheel is green and the compressor crankshaft is light magenta. The motor armature and the compressor reciprocating parts are indicated by the light orange disks on those shafts.

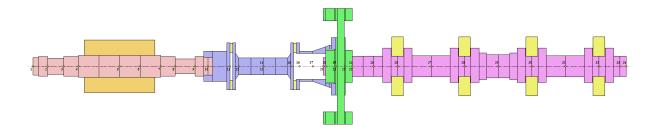
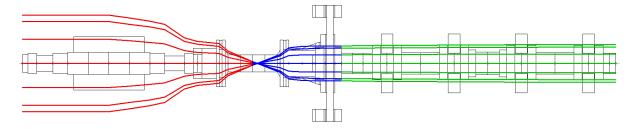


Figure 9 - Reciprocating Compressor Train Cross Section

The torsional resonance for this train has a direct torsional resonance interference with 4X compressor speed. Four times 593 RPM is the compressor excitation frequency of 2,372 CPM. Since there is significant 4X pulsation energy from the compressor, this is a resonance to be avoided.

The torsional resonance mode shape is plotted in figure 10. The horizontal axis is the axial length of the shafting. The vertical axis is the non-dimensional relative twist amplitude. This is dimensionless since there is no forcing function defined in an undamped steady state torsional



analysis.

Figure 10 - First Torsional Resonance Mode Shape - Original Coupling

In order to determine if the compressor pulsations are significant enough to cause any damage to the train components, a time transient forced response analysis is required. That means that all the alternating torque harmonics must be consolidated into an equivalent complex alternating torque waveform. This input is applied at the compressor shaft and the time based response of the shaft stresses in the train are determined.

#### **Forced Response and Shaft Stress Analysis**

The primary torsional excitations in this system are generated by the compressor. The magnitude and frequency content of these pulsations depends on the manner in which the compressor is operated. The worst case with respect to alternating torque is full load and that was the only case examined. The excitation from the compressor is expressed by the equation:

$$M(t) = Mt0 + \sum [Mt(J) * COS(J*\omega*t-\Phi(J))]$$

Where:

M(t) = Instantaneous torque at any time t, IN-LB

Mt0 = Average static torque = 54,627 IN-LB

J = Order (1 through 12)

Mt(J) = Amplitude of the J<sup>th</sup> alternating torque, IN-LB

 $\omega$  = Rotational frequency, RAD/SEC = 62.1 RAD/SEC

 $\Phi$  = Phase angle of the peak J<sup>th</sup> order torque, RAD

Table 3 shows the magnitude and phase of the first 12 orders of excitation at design load.

Order, J	Magnitude, IN-LB	Phase, Degrees	
1	8.9	192.9	
2	37,182	255.5	
3	0	156.2	
4	16,932	207.1	
5	0	215.3	
6	1,372	100.9	
7	0	163.2 329.6	
8	3,487		
9	0	167.4	
10	540	288.2	
11	0	189.9	
12	664	143.8	

Table 3 - Compressor Pulsation Data as Provided

Using this equation and order data, a value for M(t) was calculated for every degree of crankshaft rotation. One rotation at 593 RPM is 0.10118 seconds so the instantaneous amplitude was calculated for every  $2.81 \times 10^{-4}$  seconds. This waveform, plotted in figure 11, shows a period of 0.2 seconds or slightly less than 2 complete shaft rotations.

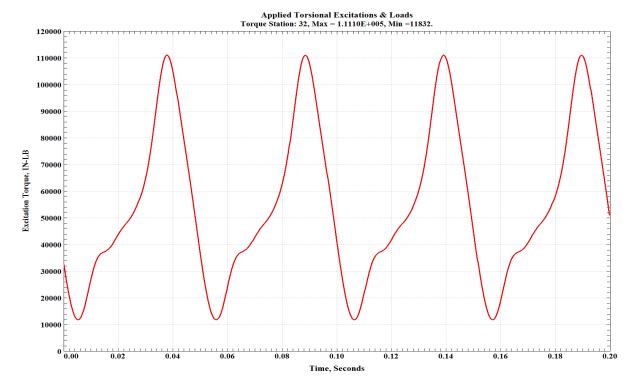


Figure 11 - Periodic Alternating Torque Function from Table 3

Since only the most highly stressed shaft section needs to be used in the comparisons, the shaft stress at resonance for an arbitrary one degree of twist angle was calculated for all shaft sections. These stress levels, shown in figure 12, are not the actual predicted shaft stresses but rather indicate that the smallest shaft diameter in the motor is the train location where the maximum torsional shear stress will occur. This is indicated by the dark red color in this area and is designated as element 8 in the finite element model.

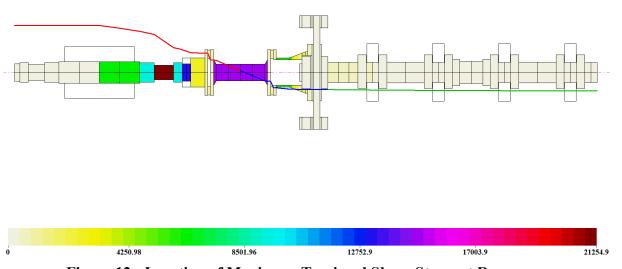


Figure 12 - Location of Maximum Torsional Shear Stress at Resonance

Figure 13 is the calculated alternating shaft torque as a function of time (0.2 seconds). Figure 14 then shows the equivalent shaft stresses in the motor shaft. These stresses are not high enough to cause long-term fatigue damage and the design is acceptable.

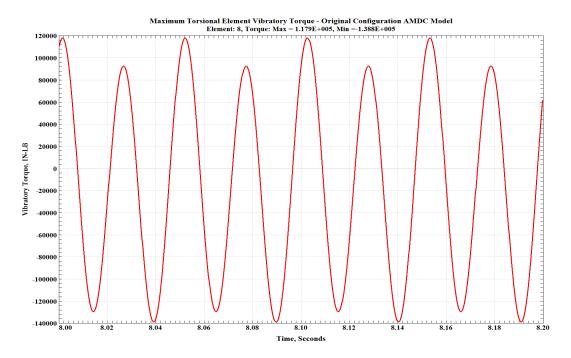


Figure 13 - Alternating Torques in Motor Shaft

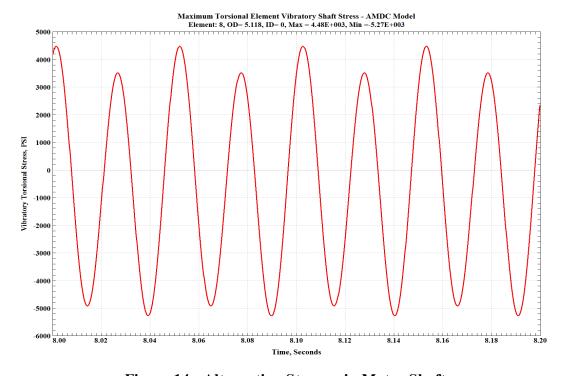


Figure 14 - Alternating Stresses in Motor Shaft

#### **Instantaneous Changes**

Time transient analysis is also required to evaluate what happens torsionally and laterally to a machine train in the event something occurs instantaneously, or at least in a very short time period. From a torsional excitation standpoint there are several events that can trigger torsional oscillations. The most serious is an internal phase-to-phase or phase-to-ground short circuit. A short circuit in the generator of a large turbine-generator train can be catastrophic. Another electrical fault that can be an issue is sudden loss of power and almost immediate re-energizing as when a breaker would open and close rapidly. A more common occurrence is a sudden load shed when the generator goes from full load to none. This does not happen instantaneously. It takes several cycles for the magnetic flux field to decay exponentially. By using this load change as the time transient input function we can calculate the transient torsional response of the entire train.

Consider the turbine generator train shown in figure 15. This is an 850MW train consisting of an HP turbine, IP Turbine, two LP Turbines, a generator and the exciter all rigidly coupled operating at 3,600 RPM. The torsional analysis was conducted to determine if a replacement generator was suitable and to investigate the effects of short circuits and load shed.



Figure 15 - 850MW Turbine-Generator Train

The shafting in this train is very large. The LP turbines have 21 inch diameter journals. The entire train is 148 feet long and the combined rotor weight is nearly half a million pounds. The smallest diameter shaft is in the exciter section and this is where the highest shear stresses will be present right after a load shed. The first two torsional resonances will comprise nearly all of the transient torsional energy. Figure 16 is a combination first torsional resonance mode shape and relative stress plot. An arbitrary one degree of twist amplitude is assumed for the resonance. The greatest twist is in the exciter shafting and where the dark red coloration indicates the highest shear stresses at the first resonance (station 76).

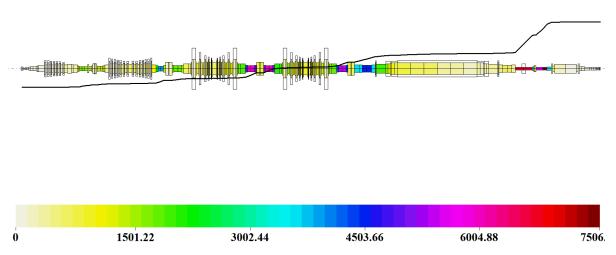


Figure 16 - Combination 1st Torsional Resonance Modes Shape and Stress Plot

The second torsional critical speed mode shape and stress plot, figure 17, indicates that the same shaft section has the greatest relative shear stress at resonance. The torsional resonances higher in frequency are difficult to excite with this type of event and are not included in this analysis.

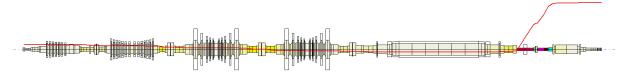




Figure 17 - Combination 2<sup>nd</sup> Torsional Resonance Modes Shape and Stress Plot

The generator manufacturer stated that the magnetic field in the generator would collapse exponentially in 60 milliseconds (3.6 cycles). The input function that was used in the analysis is shown in figure 18. The torque starts at over 20 million inch-pounds, which is normal load on the generator, and decreases to a nominal residual value after the 60 millisecond decay time.

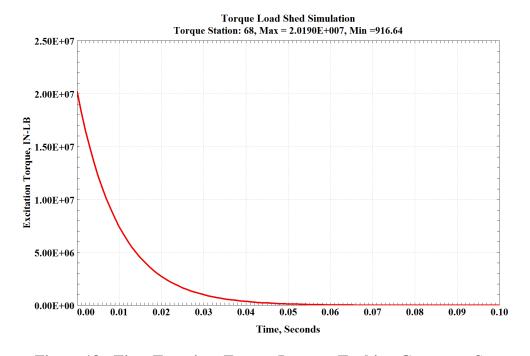


Figure 18 - Time Transient Torque Input to Turbine-Generator Set

The sudden release of torque causes the entire train of respond torsionally. All of the torsional resonances are excited but only the first two contribute significantly to the shaft stresses in the exciter shafting. An FFT of the response at this section, figure 19, shows that the first mode at 1,240 CPM dominates and the second torsional resonance at 1,566 CPM has some effect. This plot is normalized to a maximum of unity because the machine owner did not want the actual values released.

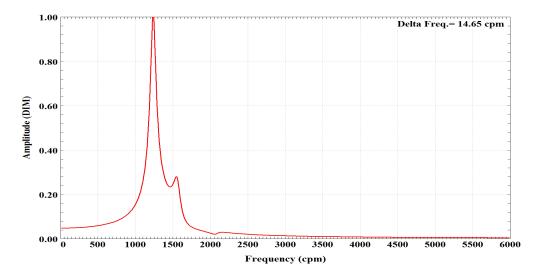


Figure 19 - Spectrum Analysis of the Transient Torsional Response

The calculated maximum shaft stresses are plotted in figure 20. This sudden torque loss causes shaft shear stress oscillations that include primarily the first torsional resonance, some second mode content and a small amount of higher mode content during the first few decay cycles.

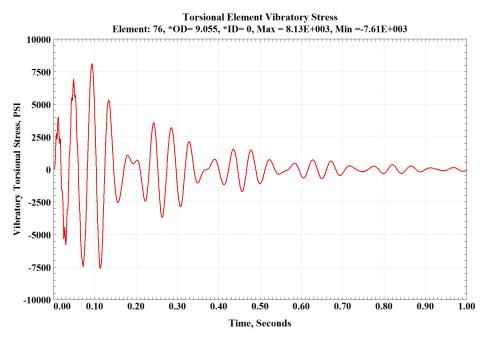


Figure 20 - Torsional Shaft Shear Stress after Load Shed

#### **Lateral Vibration Time Transient Analysis - Sudden Imbalance**

Here we will evaluate what happens to the radial vibration of this same large turbine generator train if a blade loss occurs. A blade loss is an instantaneous imbalance. While not common, this type of failure has occurred and nearly always results in a catastrophic failure and many millions of dollars of loss. In the example shown in this paper, only a portion of a blade loss will be explored to increase the clarity of what is happening. Assume that the same turbine generator set now has bearings and a random set of first and second mode imbalances indicated by the group of three hollow red arrows pointing up from each machine as shown in figure 21. The solid red arrow at station 59, the last blade row of the second low pressure turbine indicates where the sudden blade loss simulation will occur.

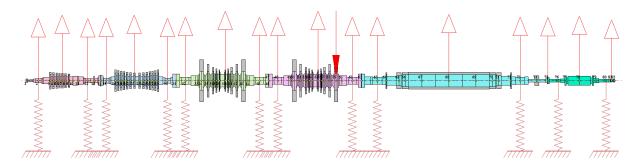


Figure 21 - Turbine Generator Set with Imbalances and Blade Loss Indication

With the train operating at 3,600 RPM, the time transient analysis is started. There is an initial transition period where the effects of the suddenly applied residual imbalances settle-out and a steady state synchronous response is established as in figure 22.

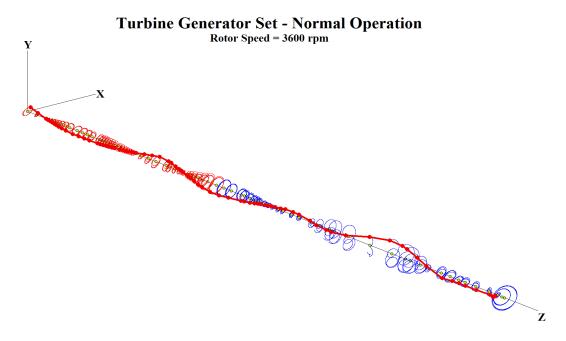


Figure 22 - Normal T-G Set Synchronous Vibration Prior to Blade Loss

It should be noted that figure 22 reveals some interesting additional information. While most of the train is showing stable round circular orbits, the generator section is not. This is because the generator is unstable and is vibrating with both 1X and subsynchronous vibration components. This is another benefit of the time transient analysis procedure.

Once the blade loss occurs at an arbitrarily chosen time of 0.1 seconds, the amplitudes increase, spiraling out to a maximum of 14.6 mils peak-to-peak before settling back to a stable 10.2 mil peak-to-peak amplitude orbit. The 3D view of this, figure 23, follows the normal operation plot previously shown. The previous smaller orbits can be seen inside the new large vibration orbits.

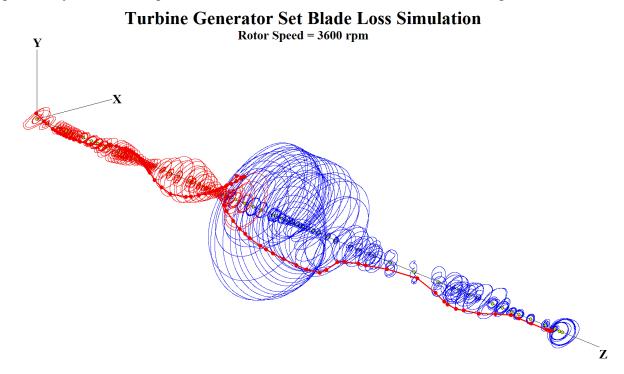


Figure 23 - Time Transient of Train Vibration after Blade Loss

As expected, the rotor that lost the blade has the highest 1X amplitude afterwards. However, the entire train is affected with most of the other rotors showing at least double the normal vibration limits. Figure 24 shows the orbit during the transition from normal operation to post blade loss operation at the station where the blade loss occurred. The black dots on this plot represent the onceper-turn timing mark. Note how the phase angle indicated by the dots abruptly changes when the blade loss occurs.

The reason that a time transient analysis is more valuable than a simple linear unbalance response analysis is that there is a period right after a transient event where the shaft vibration can be momentarily much higher in amplitude than the steady state amplitude predicted by the linear unbalance response calculations. This higher amplitude transient vibration can be enough to cause serious damage through rubs or elevated forces. The bearing forces also increase momentarily as shown in figure 25. These transient forces could be high enough to cause a bearing failure.

# X-Y Plot at Rotor Speed: 3600 rpm

Station: 59 X disp: Min = -4.4692, Max = 3.7335 V disp: Min = 3.6800, May = 6.5520

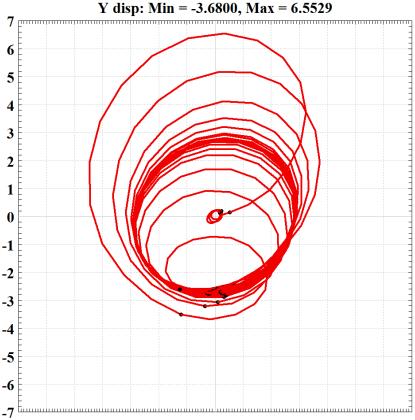


Figure 24 - LP Turbine Shaft Orbit during Blade Loss Event

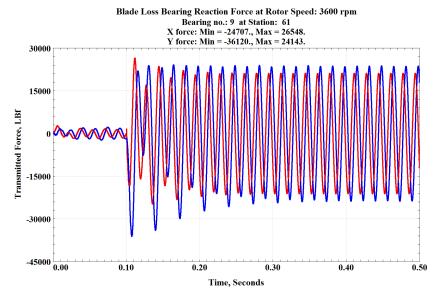


Figure 25 - LP Turbine Bearing Forces during Blade Loss Event

### **Non-Linear Bearing Effects**

Journal bearings work by supporting the weight of the rotor on a thin film of lubricant that develops a pressure profile from the act of shearing the lubricant between the rotating journal and the stationary bearing surface as indicated in figure 26. Two pressure profiles are indicated. The one on the left is a "normally loaded bearing and the one on the right is "heavily loaded.

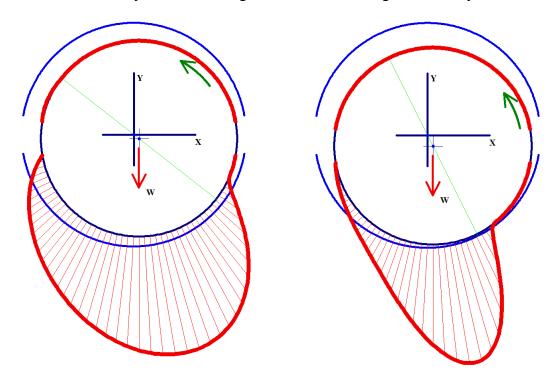


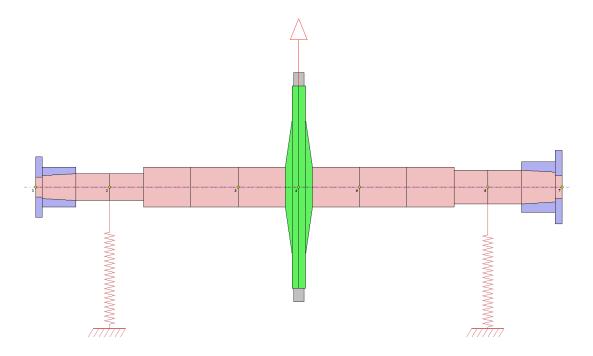
Figure 26 - Normal and Heavy Loads on a Plain Journal Bearing

The way computer programs calculate bearing stiffness and damping coefficients is to calculate an equilibrium position (eccentricity and attitude angle) of the journal in the bearing where the gravity load (plus external vector loads if present) are exactly balanced by the hydrodynamic forces indicated by the pressure profile lines in figure 26. Since the bearing stiffness is in units of Lbf/IN, the program mathematically "displaces the journal equilibrium a very small distance, X in the horizontal direction and Y in the vertical direction and then calculates the resultant hydrodynamic reaction force, F. Thus, Fx/ X and Fy/ Y become the directional principal stiffness coefficients. The cross-coupled coefficients are derived from the force that is calculated for an orthogonal displacement. The determination of the damping coefficients is done in a similar manner except, instead of a small displacement, the journal is given a small change in velocity, V. Then the Vx/ F becomes the principal horizontal damping coefficient at that speed and so forth. This method works very well as long as the displacements and velocity changes happen in an area where the hydrodynamic pressure field is fairly uniform. That is, the vibration amplitudes are relatively small compared to the distance between the eccentricity position and the bearing surface.

With a normal load situation this example bearing will operate at about 150 F at an

eccentricity of 45 percent and a maximum hydrodynamic pressure of 200 PSI. If the bearing is heavily loaded the temperature will increase to 215 F, the eccentricity to 85 percent with a peak hydrodynamic pressure of over 1800 PSI. The stiffness of the heavily loaded bearing will also be five times greater than the normally loaded case. Assume this is a 5 inch diameter journal bearing with a diametral clearance of 8 mils. This means that, under normal load, the distance between the journal center location and the bearing wall is 1.8 mils. So, a one-mil peak-to-peak orbit will comfortably fit in this space and the "small perturbation" method described above for calculating bearing coefficients will be reasonably accurate. However, if the bearing is heavily loaded, the distance between the journal surface and the bearing surface decreases to less than half a mil. Now as the shaft vibrates it will impinge on the bearing surface and, unless a rub occurs, the bearing will respond in a highly non-linear fashion. As the shaft position moves toward the bearing surface the stiffness will increase dramatically and as it moves away from the surface the stiffness will decline by as much as an order of magnitude. This severely affect the vibration response of the rotor and will usually lead to component failures. Likewise, even if the bearing loading is normal with reasonable eccentricity, if the vibration amplitude exceeds 75 percent of the bearing clearance, nonlinear effects will take over the rotordynamic response. Since the primary effect is stiffening, the rotor may see an apparent increase in the critical speed as the amplitude grows. This phenomena has been observed as a rotor passes through a critical speed it is said to "lock" onto the critical as speed increases. The increased bearing stiffness from the high amplitude is allowing the critical speed to increase and follow the rotational speed. Consider a single wheel turbine, figure 27. Figure 28 is the time-amplitude history of the center-of-shaft motion during a startup. The machine has plain journal bearings. As the critical speed is encountered near 2,500 RPM the high amplitudes lock the rotor resonance to running speed.

Figure 27 - Simple Single Wheel Turbine Rotor



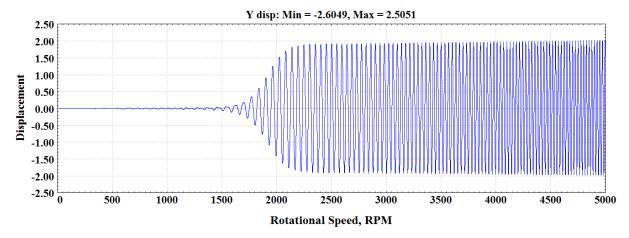


Figure 28 - Time transient Startup in Non-Linear Bearings

# **Non-Linear Effects in Squeeze-Film Dampers**

A squeeze-film damper (SFD) is a device that surrounds a regular bearing with a thin non-rotating oil film that will effectively reduce the bearing stiffness and increase damping to control critical speed amplitudes and increase rotor stability. Dampers can be applied to both anti-friction bearings and to regular journal bearings or tilting-pad bearings. During the recent design of SFD bearings for a multi-stage centrifugal compressor, the stiffness of the squeeze-film damper bearings was found to be highly dependent on the eccentricity in the damper as indicated in figure 29. The characteristics are fairly constant up to about 40 percent eccentricity when the stiffness and damping begin to increase rapidly. Figure 30 is a plot of the damping available as a function of damper eccentricity for the damper fitted to this large centrifugal compressor. Thus, significant effort was expended to design a centering device that would keep the eccentricity as low as possible when the full weight of the rotor was applied.

Typically, using a variety of methods, squeeze-film dampers are designed to be centered under normal static loads. This minimizes the direct stiffness. As static loads change or the vibration amplitude increases, the SFD direct stiffness increases. After 40 percent eccentricity the stiffness increases dramatically which is detrimental as it suppresses available damping. The damping from the squeeze-film damper is fairly constant up to the 40 percent eccentricity level when it also rises. Unfortunately, the increase in stiffness negates the higher damping values at higher eccentricities.

Figures 31 and 32 compare the runup of the steam turbine in figure 27 when squeeze-film damper bearings were applied. This plot has a relatively low imbalance of 16W/N in the center of the rotor. The first and third critical speed response peaks can be seen. When the imbalance is increased by a factor of 10, the unbalance response plot is altered to look like figure 32. The rotor displacement amplitudes have increased non-linearly, that is, by more than a factor of 10. The critical speed frequencies have also increased due to the stiffening of the dampers from the high operating eccentricities.

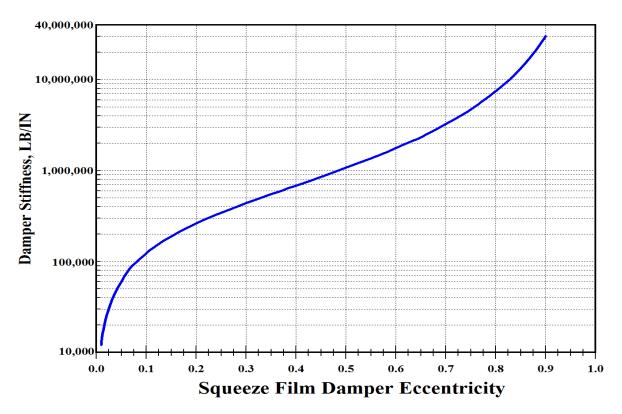


Figure 29 - Squeeze Film Damper Stiffness vs. Eccentricity

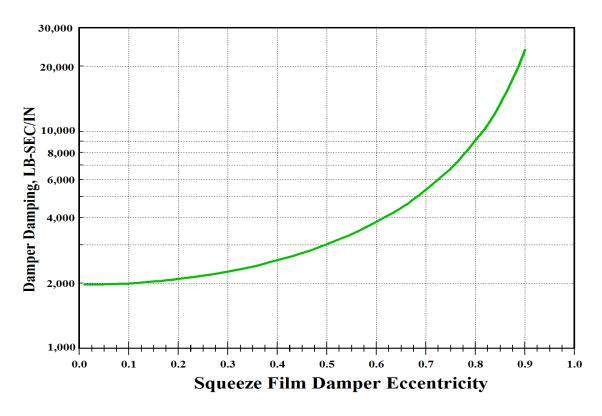


Figure 30 - Squeeze Film Damper Damping vs. Eccentricity

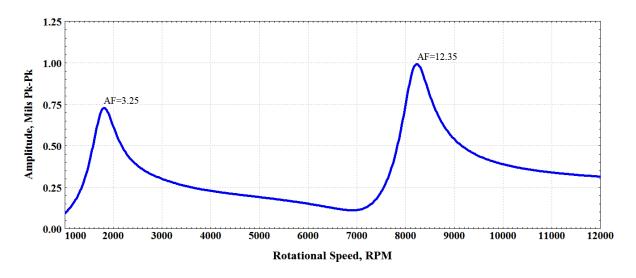


Figure 31 - Turbine Runup with Squeeze-Film Damper Bearings - Low Imbalance

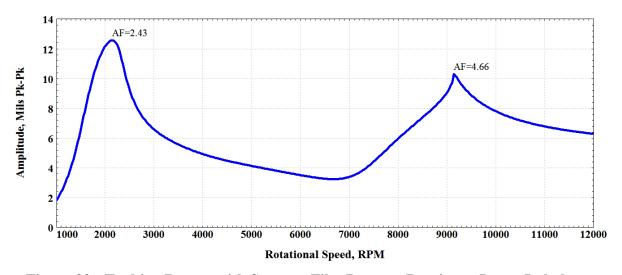


Figure 32 - Turbine Runup with Squeeze-Film Damper Bearings - Large Imbalance

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